Statistical Physics

Abigail Tan January 18, 2023

The Microcanonical Ensemble

This represents the possible states of an isolated system whose total energy is exactly specified.

Number of states. $\Omega(E)$ = number of states with energy E; have $p(n) = \begin{cases} \frac{1}{\Omega(E)} & \text{if energy} \\ 0 & \text{otherwise} \end{cases}$

Entropy. $S(E) = k_b \log \Omega(E)$ (additive)

Temperature. $\frac{1}{T} = \frac{\partial S}{\partial E}$ Heat capacity. $C = \frac{\partial E}{\partial T}$ and $\Delta S = \int_{T_1}^{T_2} \frac{C(T)}{T} dT$ Pressure. $p = T \left(\frac{\partial S}{\partial V}\right)_E$

First Law of Thermodynamics. dE = TdS - pdV

Heat capacity (in terms of S). $C_V = T\left(\frac{\partial S}{\partial T}\right)_V$ and $C_p = T\left(\frac{\partial S}{\partial T}\right)_p$

The Canonical Ensemble

This represents the possible states of a system in thermal equilibrium with a heat reservoir at fixed temperature.

Approximation for $\Omega(E_{\text{tot}})$. $\Omega(E_{\text{tot}}) \approx e^{S_R(E_{\text{tot}})/k_B} \sum_n e^{-E_n/k_B T}$ Boltzmann distribution for CE. $p(n) = \frac{e^{-E_n/k_B T}}{\sum_m e^{-E_m/k_B T}}$ Partition function. $Z = \sum_n e^{-\beta E_n} \implies p(n) = \frac{e^{-\beta E_n}}{Z}$ where $\beta = 1/k_B T$ Energy statistics for CE. $\langle E \rangle = -\frac{\partial}{\partial \beta} \log Z$, $\Delta E^2 = -\frac{\partial}{\partial \beta} \langle E \rangle = k_B T^2 C_V$ CE entropy. $S = -k_B \sum_n p(n) \log p(n)$ and $S = k_B \frac{\partial}{\partial T} (T \log Z)$ Free energy. F = E - TS, $F = -k_B T \log Z$ (most likely E minimises F)